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THREE STAGE ATTENUATION USING QUANTUM PLASMA PRECIPITATION

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Figure

THE BOSTONIAN, NOVEMBER 18, 1861.

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PHOTO - PCS - MARCH 1981

**Spouse and children - "or anything therefor
the spouse will take, if necessary
and no cause shall ever exist."**

PL 76-40726

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and subsequent to being advised the chairman stated much PL-42 "further
work will be done" it was decided to keep another a month meeting the

The author's report was taken subject to a reference to multiplication

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REPORT DOCUMENTATION PAGE

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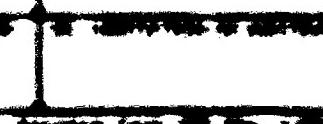
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12. ABSTRACT

Conventional three-stage attenuators were developed in the 1960's and had attenuation limited to 1.0% transmission and 0.01% insertion loss. Current development has at least two basic new concepts which can be attained using two new types of attenuators which accomplish either a reduction of a quarter wave retardation added to transmission with low insertion loss. The purpose of the current effort was to test the basic concepts using commercially made plastic optical fiber. Various plastic materials provide attenuated stages and an experimental setup designed to provide various transmission for high attenuation when to determine the feasibility of using a three stage plastic attenuator with low losses in the 0.01 to 0.001 range. The basic could be used to reduce input power requirements while maintaining a constant output level. The experimental results indicate that the attenuator could only provide attenuation of 1 to 3 percent. The attenuators currently used in most all wave guides cannot be attenuated over 0.01 percent without a significant increase in size to achieve desired levels. There is a possibility that the new attenuators will be able to reduce the size of the entire system.

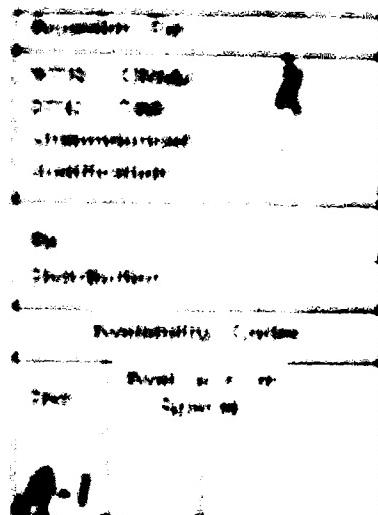
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1.0 INTRODUCTION

Conventional polarizer attenuators have been used since at least the 1920s. The transmittance of the two-stage polarizer attenuator, where one polarizer is fixed and the other one is rotated, is known as Matus' law:

$$T(\delta) / T(0) = (\cos \delta)^2$$

where $T(\delta)$ is the transmittance of the two polarizers and δ is the angle enclosed by the principal transmittance axes. This device was unreliable if the source was partially polarized or the sensitivity of the detectors varied with the angle of polarization (Ref. 1). Dowell developed the three-stage polarizer attenuator which overcame the defects in the two-stage attenuator (Ref. 2). In Dowell's method the first and last polarizers were stationary with their optic axes parallel and the middle polarizer was rotated. The three-stage polarizer attenuator transmitted an intensity governed by a \cos^4 relationship,

$$T(\delta) / T(0) = (\cos \delta)^4$$

To obtain accurate measurements, the extinction ratios of the polarizers and the birefringence of the middle polarizer must be known and accounted for when the attenuator is calibrated. In general, the conventional three-polarizer attenuators utilizing film polarizers were limited to 0.001 transmittance units (Ref. 1).

Several different types of polarizers have been used in the three-stage attenuators. Bennett used sheet Polaroid mounted between distortionless glass plates and determined photometric linearity to better than 0.1 percent (Ref. 3). Mietenz and Eckert discussed systematic errors due to imperfections in sheet polarizers, setting and alignment errors, and incident beam incidence angle and polarization. They concluded that the accuracy of three-polarizer film attenuators is limited to 0.001 transmittance unit largely because of the unknown birefringence of the middle polarizer. Instead they employed either a half-wave or a quarter-wave retardation plate with a precisely known birefringence and two sheet polarizers and obtained at least 10 times more accuracy than the conventional three-polarizer attenuators (Ref. 1).

Polarization prisms would also avoid the birefringence problem. Mietenz and Eckert did not pursue the use of prisms because of potentially serious systematic errors caused by their small field angles and because the accurate measurement of the high extinction ratios of good polarizers was difficult (Ref. 1). Bennett tested, but did not use, a good Glen-Thompson prism that deviated the beam by less than 1 min of arc.

and was a high schlieren quality calcite, because the intensity variation was not symmetric in the four quadrants. He suggested that a prism-type polarizer with adequate performance would have the advantage of a wider, more useful spectral range than was possible with sheet polarizers (Ref. 3).

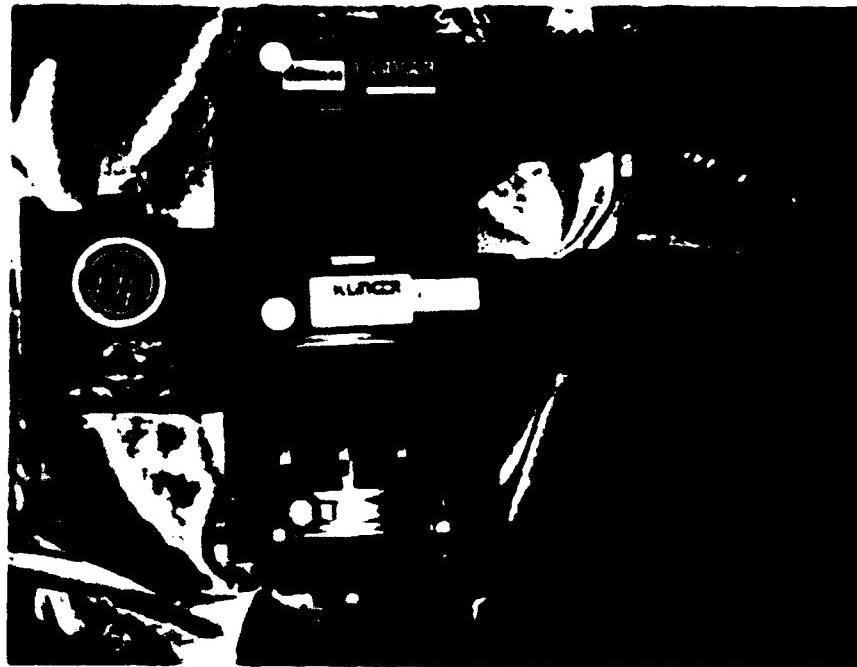
The current work was based on the same concept and used specially selected high quality optical Glan-Thompson prisms, extremely precise automated stages, and a combination of optical density filters with a lock-in amplifier, to obtain accurate measurements and to determine the feasibility of using a three-stage polarizer attenuator with laser beams in and near the visible range. This device should be effective for independently calibrating experimental signals and transmittance of neutral density filters over a wavelength range from 350 to 2500 nm and over an optical density range of nine orders of magnitude.

2.0 EXPERIMENTAL ARRANGEMENT

2.1 ATTENUATOR DESIGN

The three-stage attenuator was comprised of specially selected high quality optical Glan-Thompson prisms manufactured by Karl Lambrecht Corporation (MGT25E10-90), Chicago, Illinois. The calcite prisms are glued at the interface. The optical glue limited the capability of the attenuator at the ultraviolet end of the spectral range, but it also provided better maximum attenuation. The prisms were selected so that one crossed pair would have an attenuation equal to or greater than $10^6:1$. The prisms had a wavelength range of approximately 350 to 2500 nm and a tolerance on maximum beam deviation of approximately 6 arc sec.

The prisms were each mounted in precision stages. Two of the stages, manufactured by Klinger Scientific, Garden City, New York, were electronically driven. The first stage was required to accurately return to the 0-deg position and to have a repeatable 90-deg rotation. These criteria were easily met with one of the Klinger stages. The requirements for the second stage (middle polarizer) were the most stringent since this stage was the only one which would rotate after the attenuator was aligned. This stage was responsible for determining the angle " β " with great accuracy. Consequently, the second stage consisted of a Klinger Scientific, Garden City, New York, stage and an encoder. The encoder on the middle stage was accurate to approximately 1 arc sec. The third stage was a manual stage which incorporated both fine and coarse adjustments. Figures 1a, 1b, and 1c show the assembled three-stage polarizer attenuator taken from the top (Fig. 1a), from the front (Fig. 1b), and from the side (Fig. 1c).



(a). Top view.

Figure 1. Three-stage polarizer attenuator.



(b). Side view.



(c). Front view.

Figure 1. Concluded.

2.2 ALIGNMENT OF THE ATTENUATOR

The alignment for the stages was accomplished mechanically by using an optical post which was the same diameter as the prisms. The post was inserted through the mounts (where the prisms would be mounted) and then the mounts were adjusted and tightened down. The post was removed.

Each prism was then individually mounted in each stage position and tested. A target was placed 6 to 8 ft away from the attenuator and then the stage was rotated through 360 deg to ensure that the laser spot remained in the same location on the target. All three prisms performed properly.

An alignment procedure was developed which ensured the first and third polarizers had their optic axes in parallel positions and the second polarizer reached a minimum transmittance when it was crossed with the other two polarizers. The procedures used are described below and depicted in Figure 2. (Note: Prism 1 is mounted closest to the laser source, prism 2 is mounted in the center of the attenuator, and prism 3 is mounted closest to the detector.)

Step 1. Set all of the mounts to their 0-deg position and insert all the prisms in the same approximate orientation.

Step 2. Set prism 3 near the 90-deg or null position.

Step 3. Set prism 2 so that prisms 1 and 2 are exactly 90 deg off (cross polarized). This location should provide an exact null and prism 2 will be in approximately the 90-deg position.

Step 4. Set prism 1 in the same position as prism 2.

Step 5. Return prism 3 to the 0-deg position and adjust it until an exact null is reached (prisms 2 and 3 are crossed polarizers). Lock prism 3 in place.

Step 6. Return prism 1 to its exact original position. At this point, prisms 1 and 2 should be exactly 90 deg offset and prisms 2 and 3 should also be exactly 90 deg offset.

Step 7. Return prism 2 to its 0-deg position. The attenuator should now be set for maximum transmission.

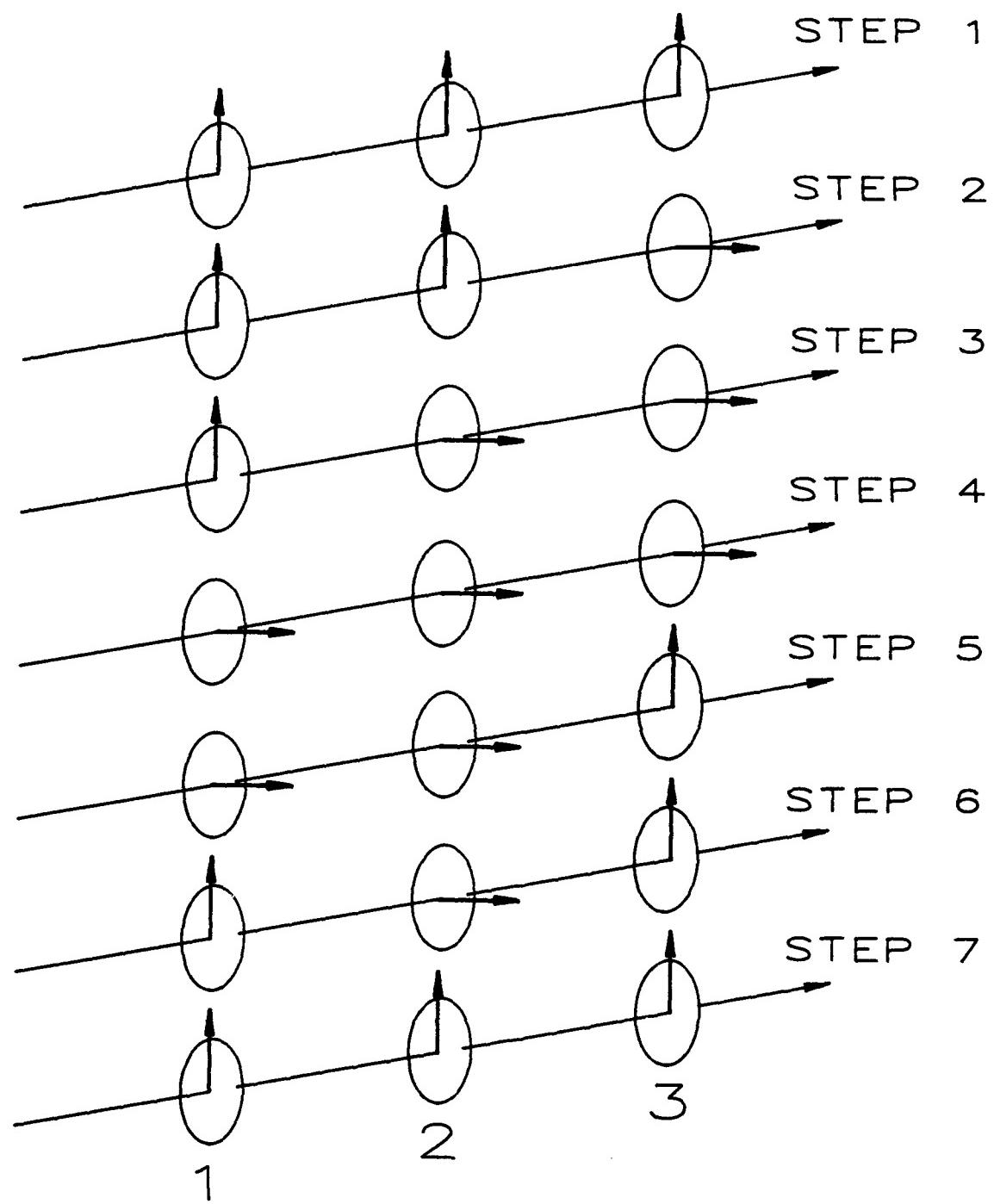


Figure 2. Attenuation alignment procedures.

2.3 EXPERIMENTAL LAYOUT

The three-stage attenuator was then tested at experimental conditions to determine range and accuracy. The experimental schematic is shown in Figure 1.

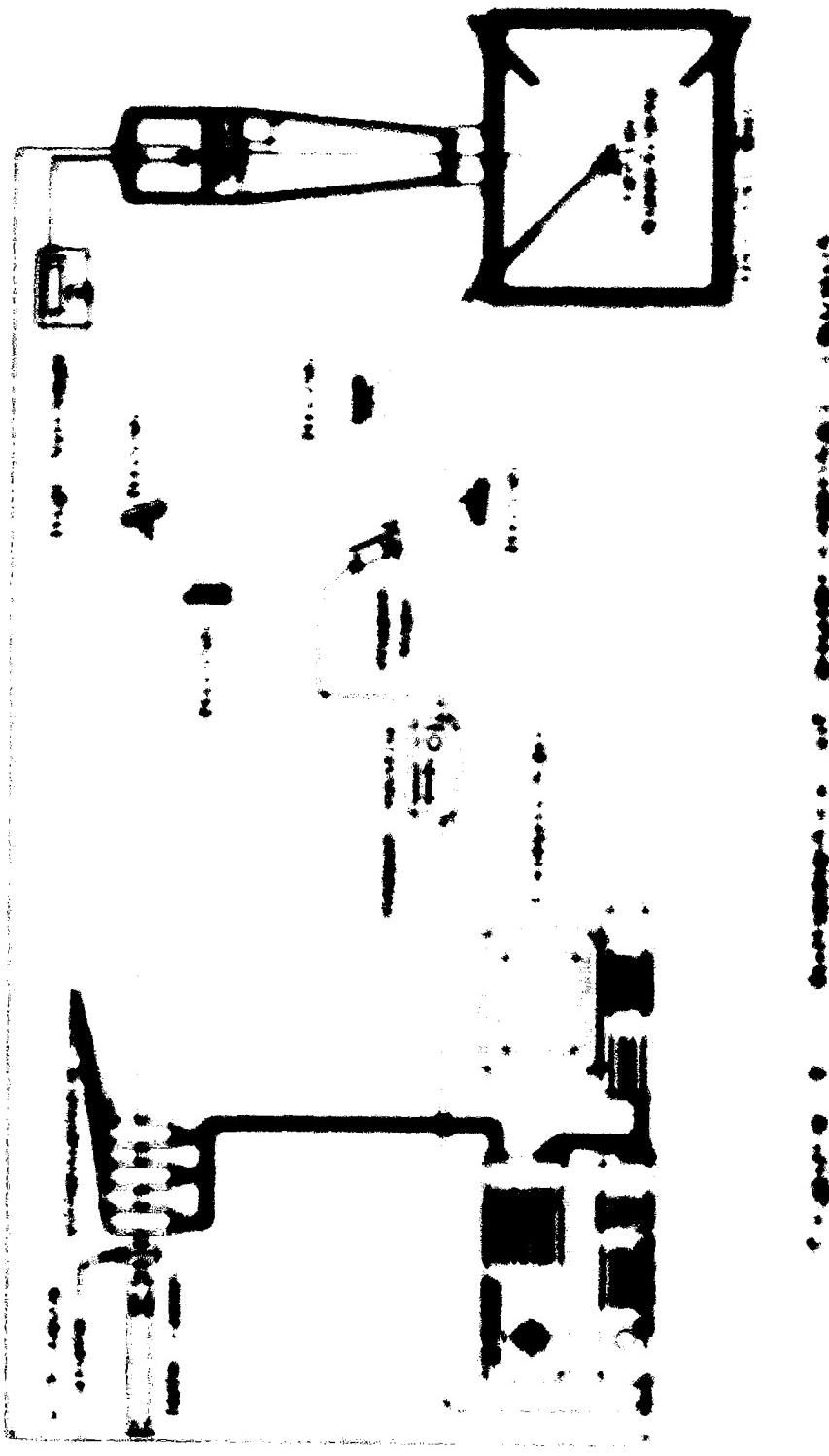
The laser was a 5-mW helium-neon laser (G3) with an adjustable beam expander. A quarter-wave plate was positioned between the laser and the first attenuator. Consequently, a monochromatic, collimated, unpolarized light source was normally incident upon the attenuator. A series of lenses were used to properly position the light for the detector.

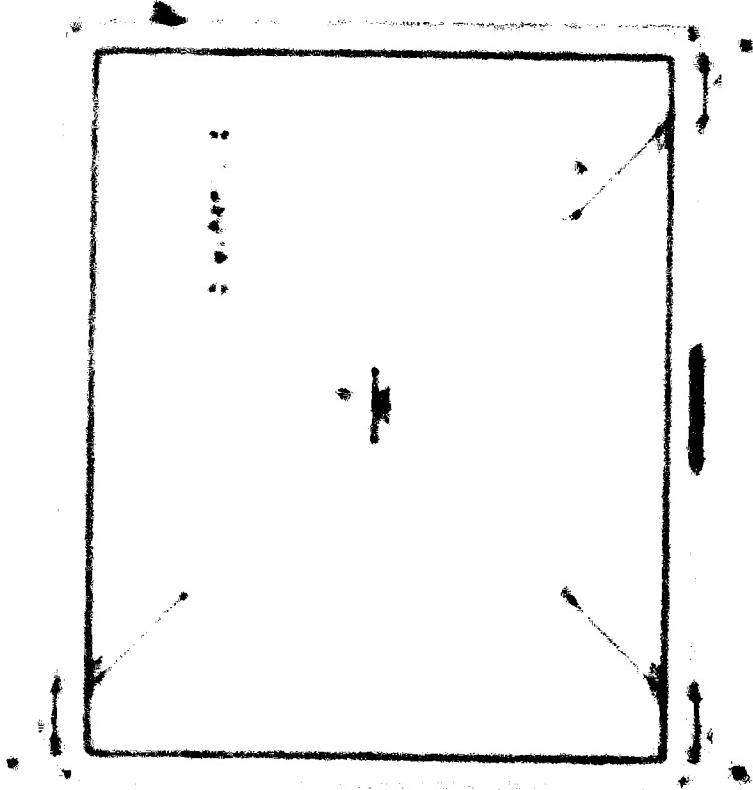
A chopper with reflective blades was used to turn part of the beam onto an integrating sphere which was coupled to a silicon photodiode. The detector saw the input laser signal and was read on the lock-in amplifier.

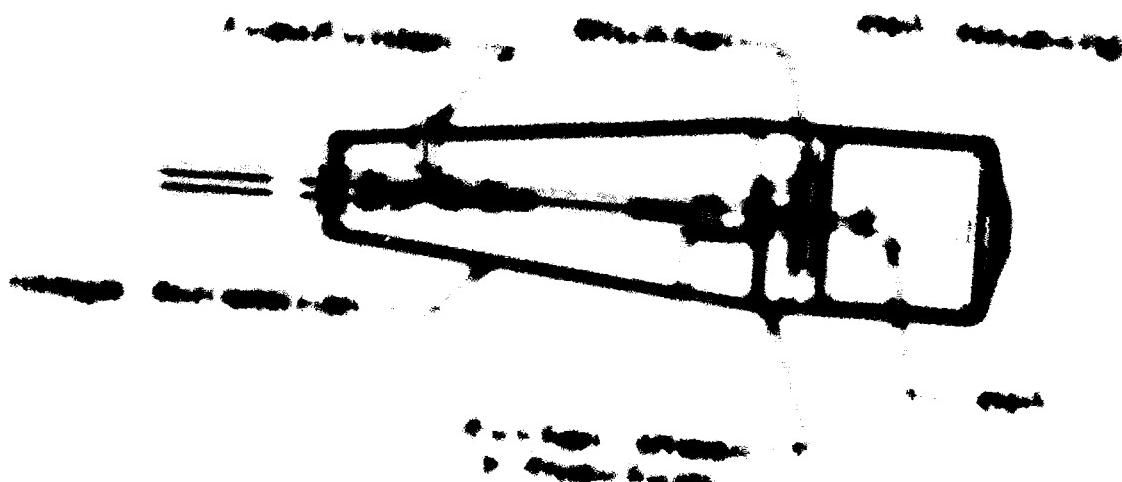
The laser light then entered a utility box measuring 33 x 33 x 48 cm. There were holes in its walls for both the incident and reflected laser light sources. The incident beam caused the laser light to scatter, in this case a polytetrafluoroethylene (Teflon) reflector (Labsphere, Inc., North Sutton, New Hampshire) was held about 10 cm from the center of the box. The laser beam entered the box at a 45° angle to the sample's normal. A portion of the reflected beam entered through a tube to the photometer. For alignment purposes, a thin tube was inserted in the box so the beam could exit when no sample was in place. Two thin tubes, each ending in a hole in the inner wall, were inserted through the small holes in the box. Through these tubes the beam entered and exited the box. Alignment of the tubes was done by eye. The laser beam was adjusted until it appeared to be centered when viewed through each of the tubes. Figure 4 illustrates the basic setup of the utility box.

The photometer, filter wheel assembly, and photomultiplier tube (PMT) have several capabilities: (1) the ability to scramble the signal source, (2) the ability to focus the beam, (3) the ability to change aperture sizes, (4) the ability to eliminate most sources of stray light, and (5) the ability to reduce noise. Neutral density filters in the optical train often the light from being focused on the optical axis. Figure 5 provides a ray diagram of this setup. The photometer has a lead housing (Products for Research, Inc., Danvers, Massachusetts) designed to provide radiofrequency shielding, a double window and a filter to reduce noise and the reduction of dark current.

The PMT was connected to a phase lock amplifier. The gain control on the photometer was used to obtain data over a six order of magnitude range. Six orders of magnitude were obtained by attenuating the high end of the signal with a National Institute of Science and Technology (NIST) neutral density filter. Consequently, even the weak signals were above the detection limitation and within the linear range of the photomultiplier.







to make a significant contribution to the development of a new model of

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the same as the one in the first sentence. The second sentence is also true, because it is the same as the first sentence.

Figure 1 shows the experimental and theoretical trajectories taken against the prior
knowledge graph as expected. The algorithm takes one iteration at 0.01 sec and at 0.5%
they converge to stable solution and remains. In addition the experimental data only
shows a single 4 components fit within a confidence. The experimental data
against uncertainty and with the theoretical data through a series of iterations converges
significantly from the theoretical data as shown. It shows a confidence with the
experimental uncertainty of the measured data points.

"...that's where the idea of the community agrees against the traditional interpretation to the idea that a day is a unit and it is day and night being where the natural community occurs. In all these countries the communities have developed approaches from the traditional idea either a collection of communities and areas of migration. The areas are at times a part of a region or small clusters. The other approach from Germany clearly also a collection of areas within a region. Consequently no extensive exchange of the people between the different regions and areas of migration.

There is one other important difference between the two models. The first model is based on the assumption that the economic environment is stable and predictable, while the second model is based on the assumption that the economic environment is dynamic and unpredictable.

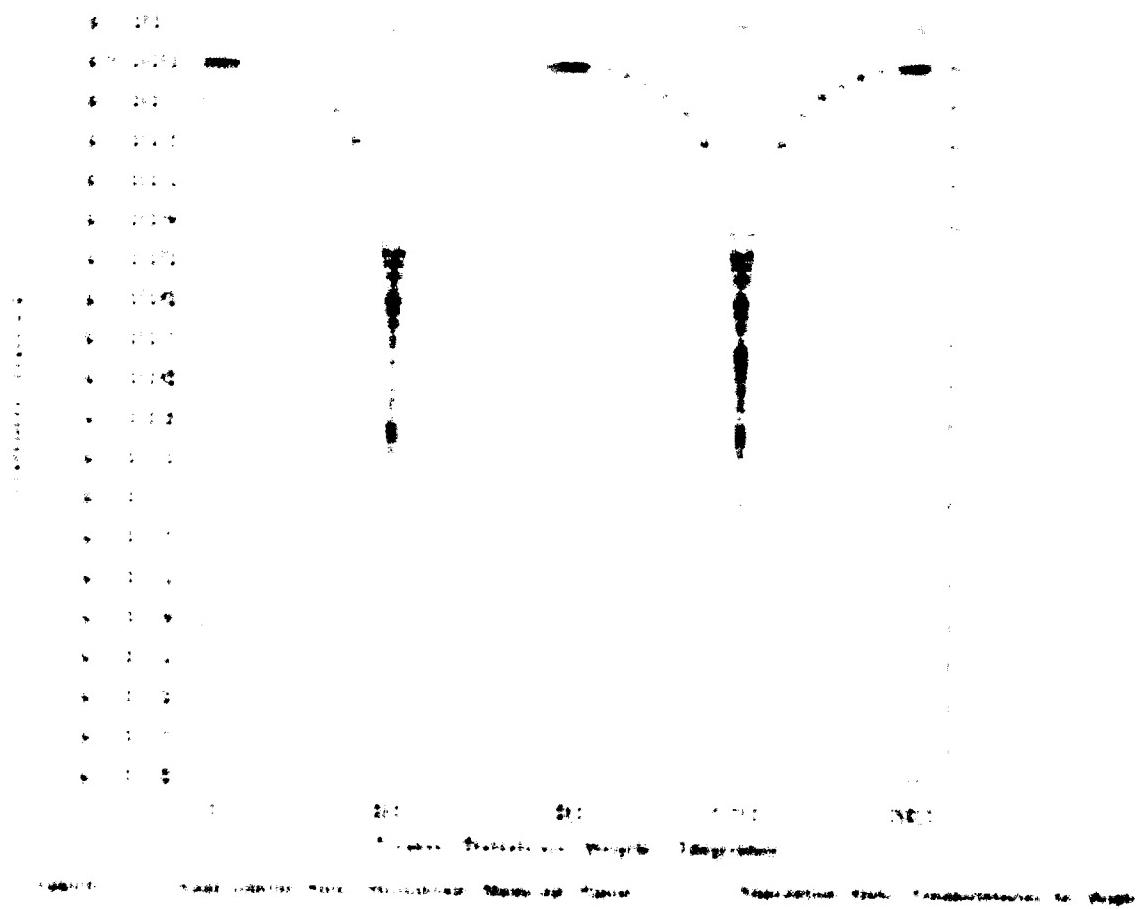


Figure 4. Comparison of the measured temperatures versus the ground truth.

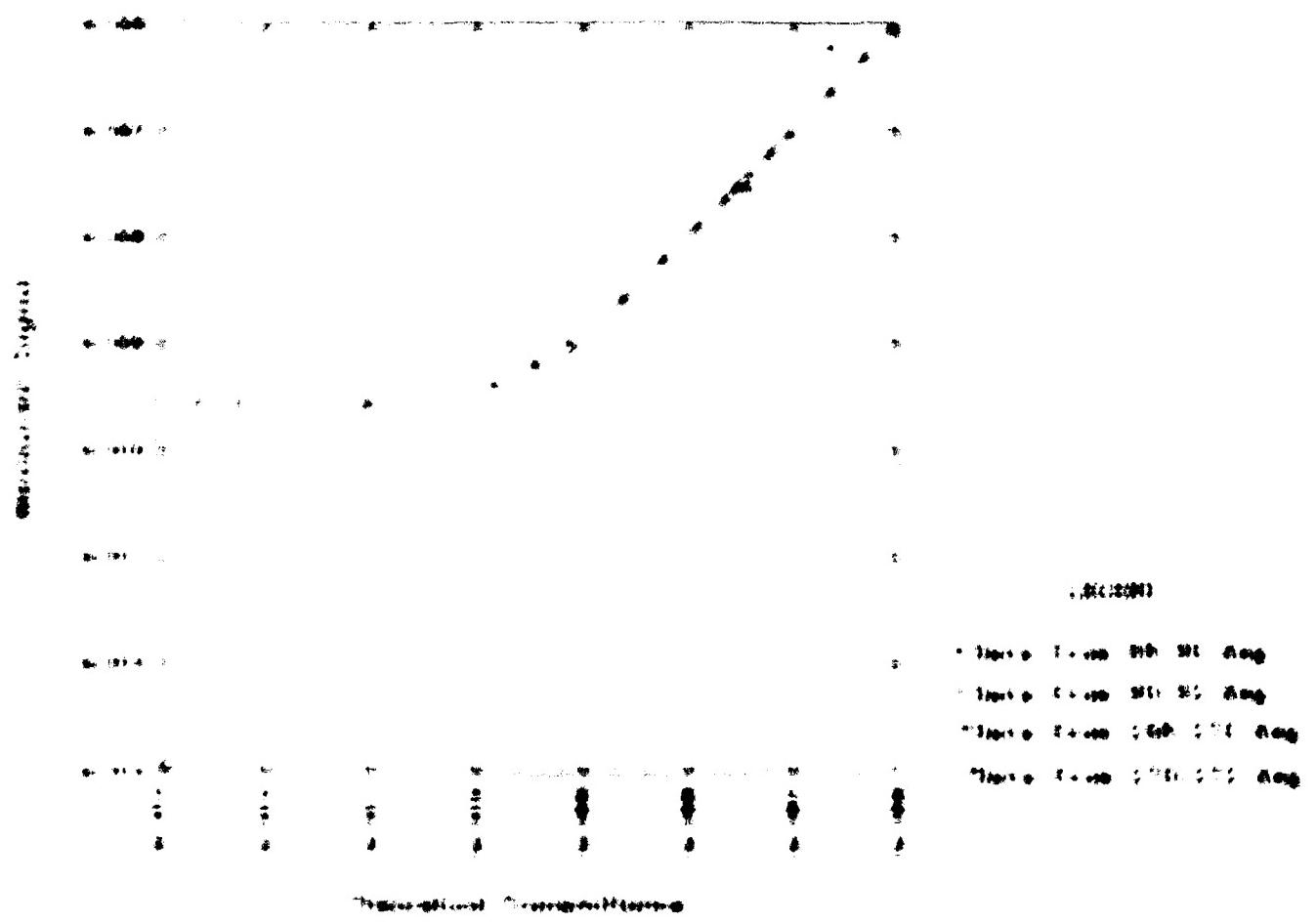


Figure 4 Measured signal versus theoretical transmission for data in the vicinity of 60 and 570 Ang

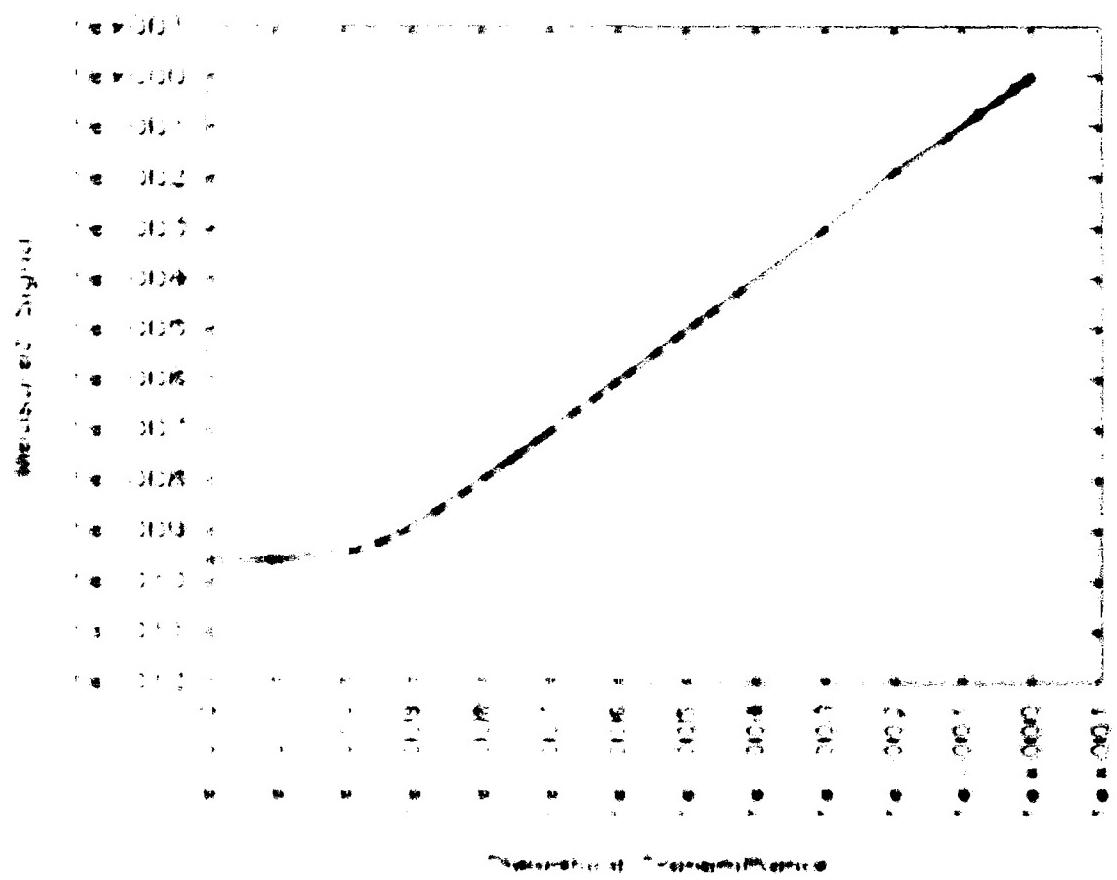
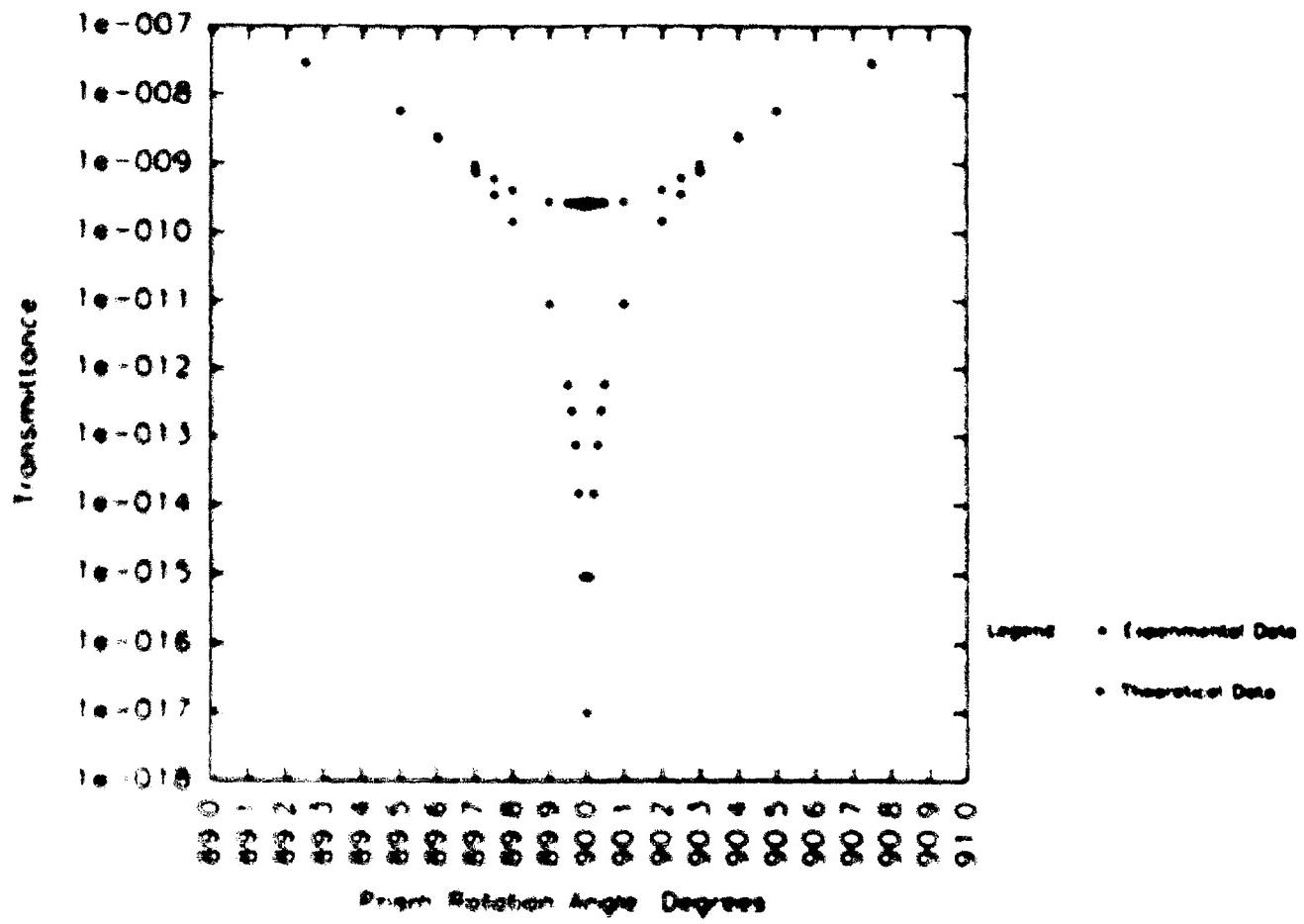
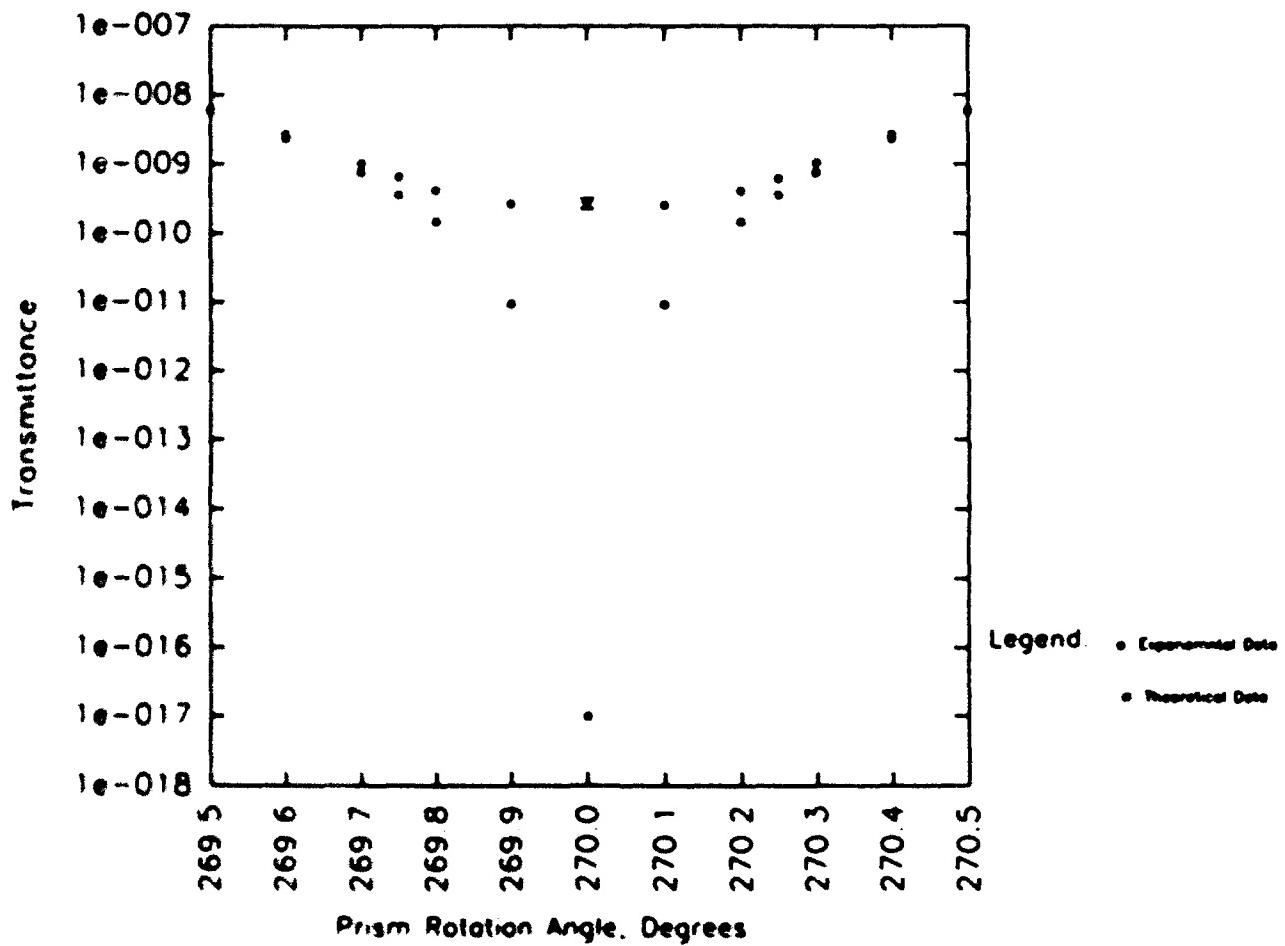


Figure 6. Measured signal versus Received Power for the entire data set.



(a) Angles near 90-deg

Figure 9 Data plot of the transmittance versus the prism rotation angle.



(b) Angles near 270 deg.

Figure 9. Concluded.

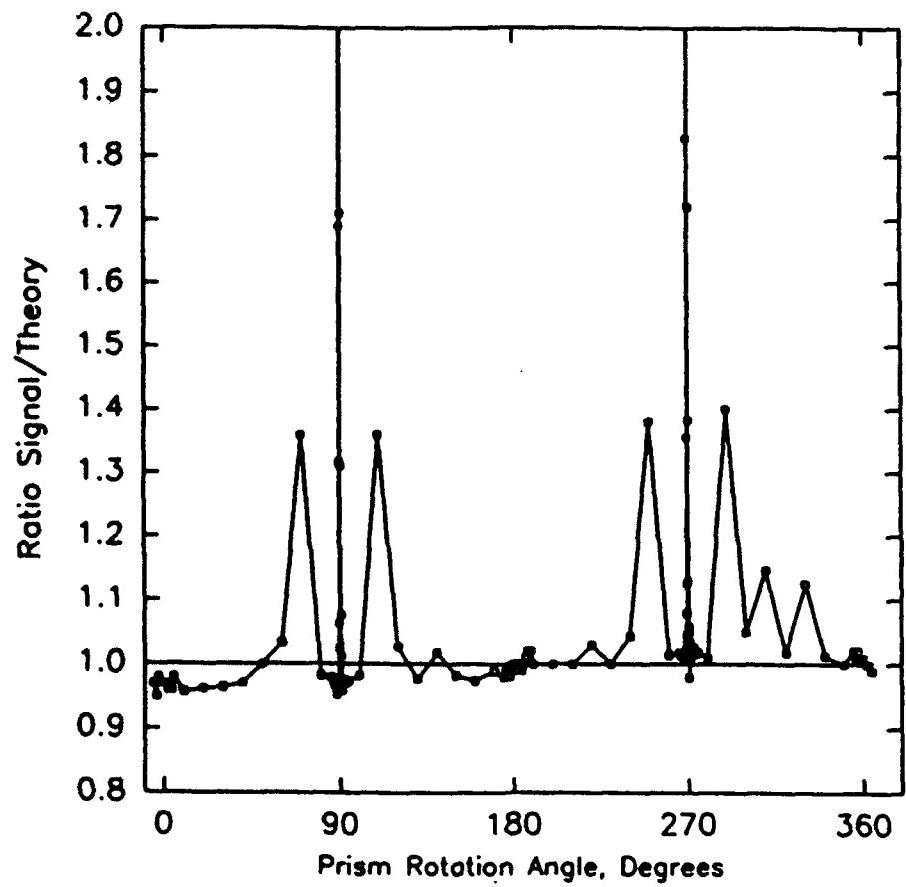


Figure 10. Ratio of the experimental transmittance to the theoretical transmittance versus the prism rotation angle.

3.2 EXPERIMENTAL ERROR

3.2.1 Measurement Error

There were several sources of measurement and random error present throughout the experiment. These errors include nonlinearity in the phase lock amplifier, noise from the photomultiplier, stray light in the optical train, variation in the lock-in amplifier reading, and drift in laser power. No attempt was made to control the laser power.

3.2.2 Systematic Error

Data collected around the 0-, 90-, 180-, and 270-deg angles indicated significant systematic errors. The major errors include drift in the laser signal from the beginning to the end of the data run and slight alignment errors in the prisms. The drift was noticed in the signal from the beginning to the end of the data runs. In addition, the equipment was left for several 30-min intervals without altering the equipment setting and experimental data did change slightly with time. Alignment errors are noticed in the lack of perfect symmetry in data around the 0- and 360-deg points. In addition, there was a small double hump pattern around the 90- and 270-deg points. Additional systematic errors would include the temperature change in the room affecting the laser detector or the electronics and small inaccuracies in the mechanized center stage.

4.0 CONCLUSIONS

The three-stage polarizer attenuator, with the advanced Glan-Thompson optical prisms, the extremely precise automated stages, and an experimental setup which allowed accurate measurements of high extinction ratios, was capable of providing attenuation of a laser beam for nine orders of magnitude with uncertainties of 1 to 2 percent. This attenuator was useful in providing practical accuracies of a few percent over a wide dynamic range with a simple technique. The data were symmetric throughout all four quadrants. The attenuator can be used to calibrate neutral density filters over a wavelength range from 350 to 2500 nm and over an optical density range of nine orders of magnitude. These results were obtained with laser sources only, and no corrections were made for drift in the laser power or signal processing.

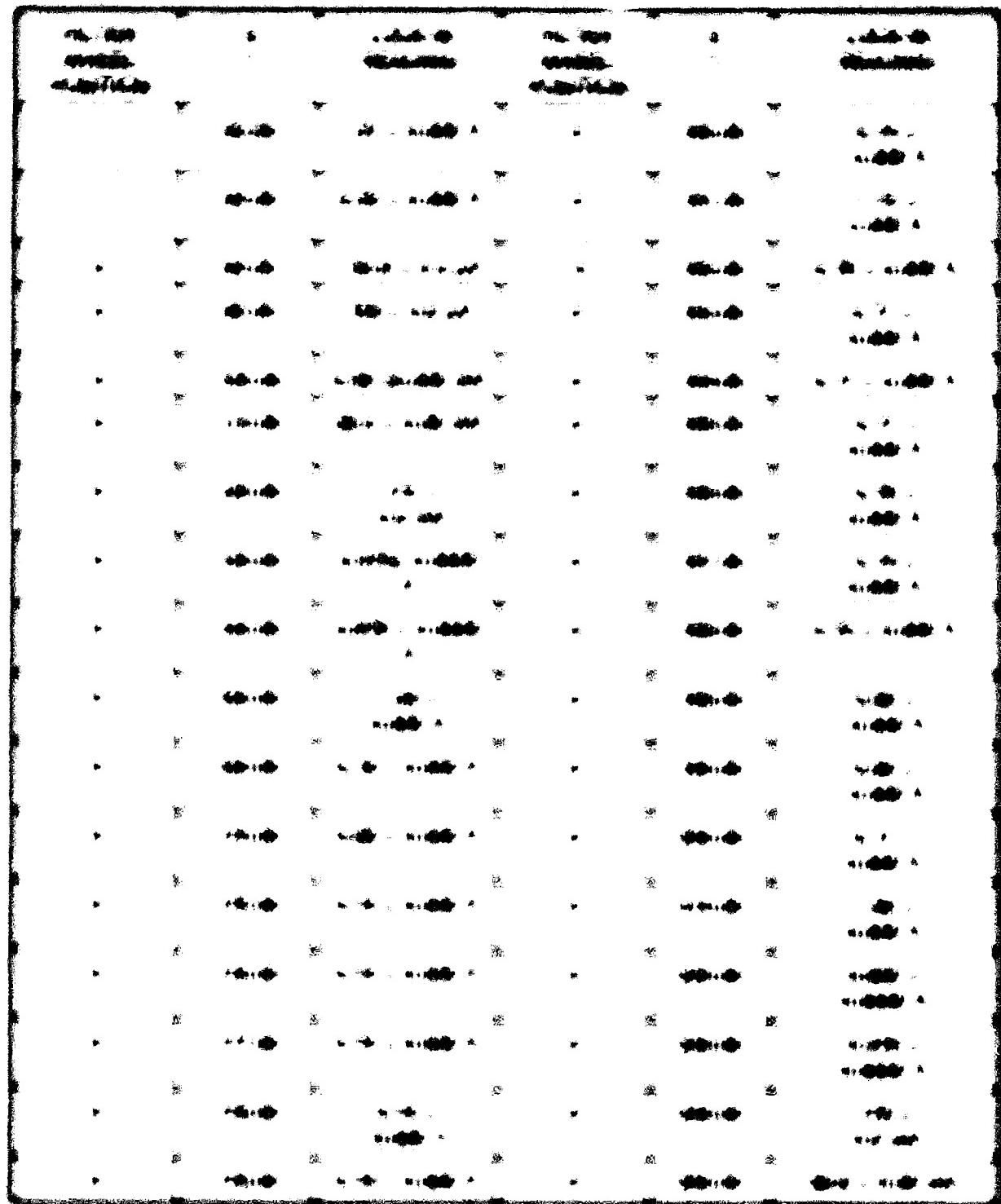
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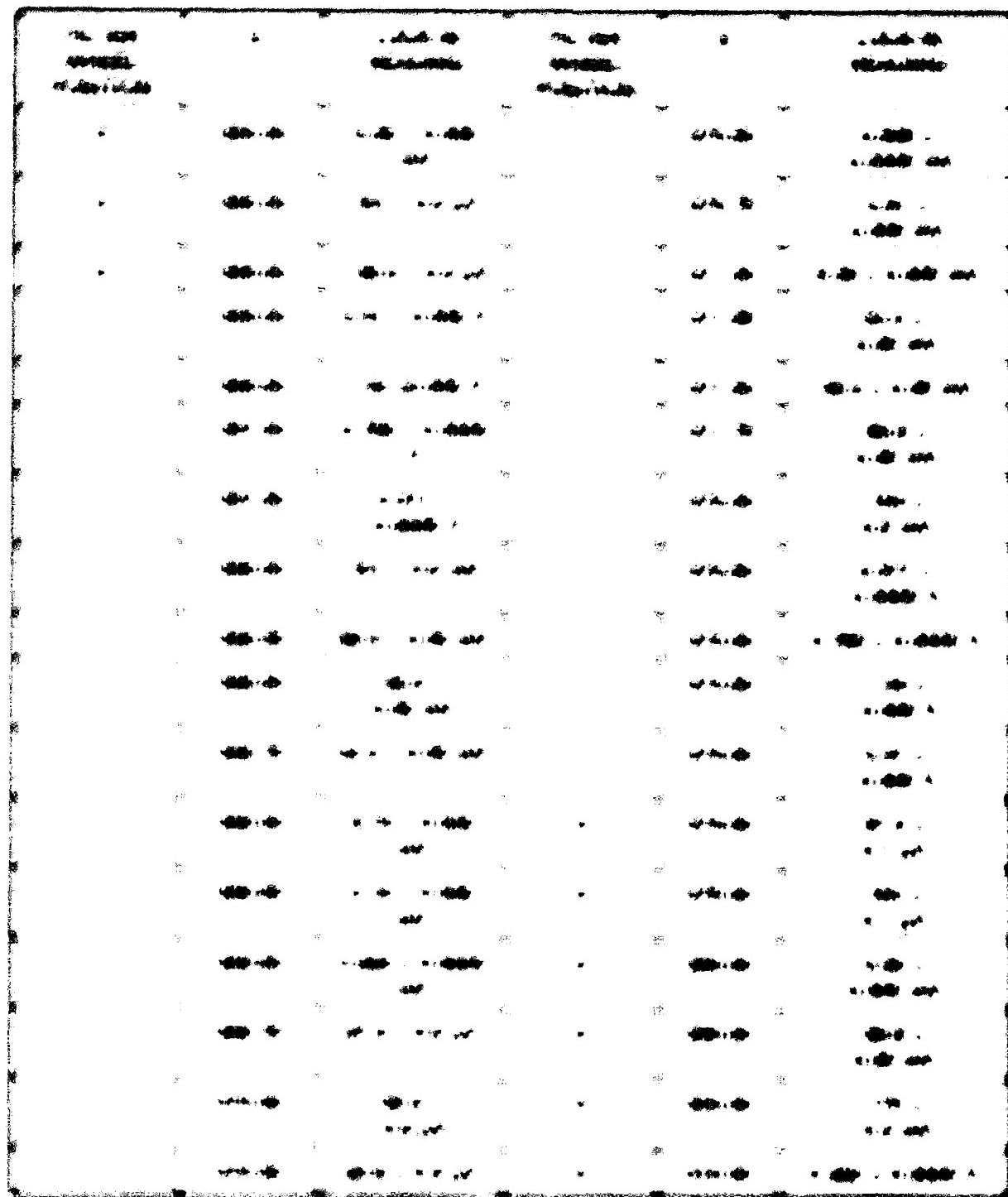
1. Mielenz, K. D. and Echerle, K. L., "Accuracy of Polarization Attenuators," Applied Optics, Vol. 11, No. 3, pp. 594-603, March 1972.
2. Dowell, J. H., J Sci Instrum, Vol. 8, p. 382, 1931.
3. Bennett, H.E., "Accurate Method for Determining Photometric Linearity," Applied Optics, Vol. 5, No. 8, August 1966.

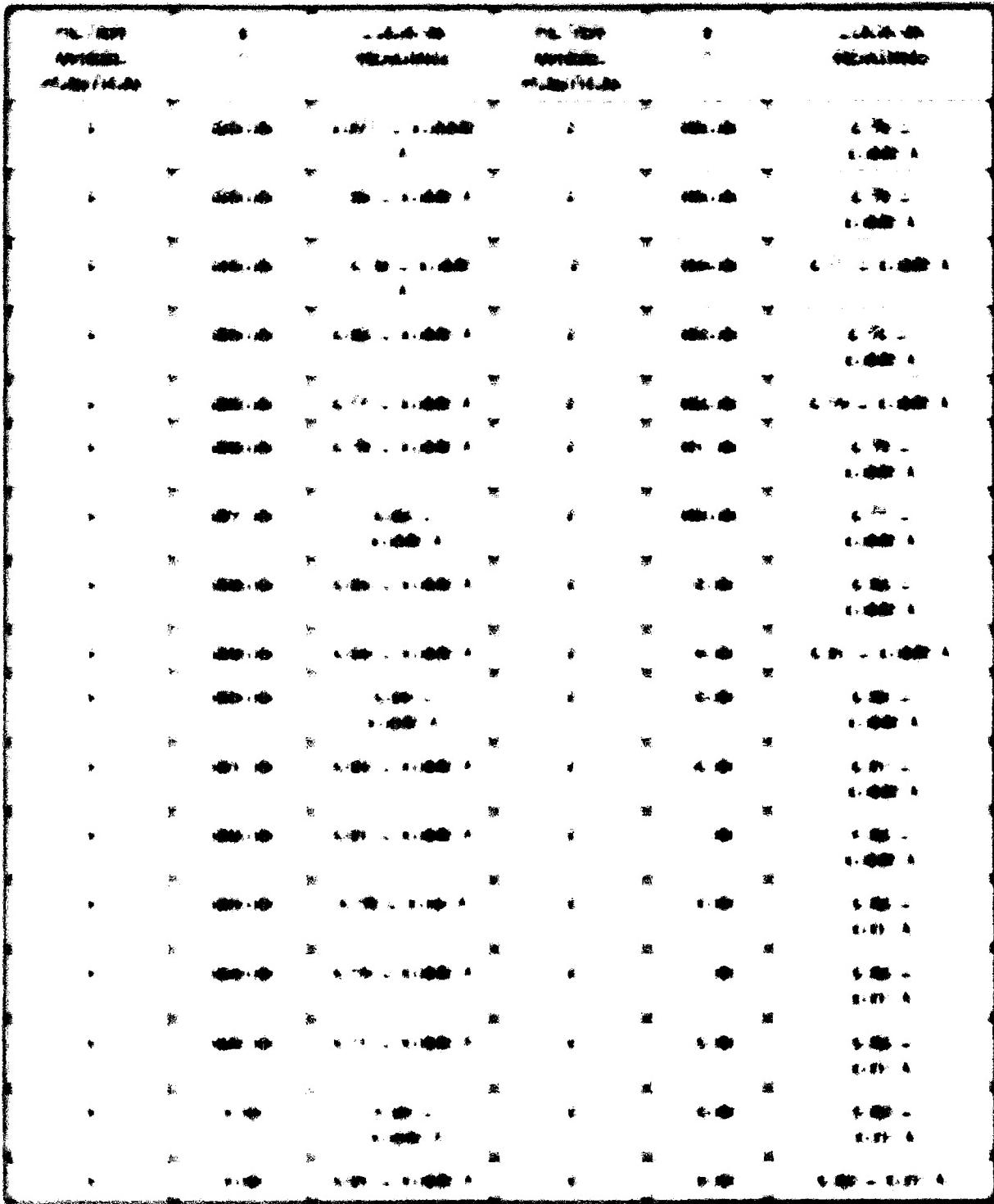
APPENDIX A

RAW DATA COLLECTION

FILTER WHEEL POSITION	0	100	200	300	400	500	600	700	800	900
0	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
1	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
2	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
3	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
4	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
5	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
6	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
7	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
8	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
9	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
10	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
11	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
12	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
13	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
14	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
15	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
16	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
17	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
18	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
19	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
20	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
21	0.00	0.00	0.							







FILTER WHEEL POSITION	θ (°)	LOCK-IN READING	FILTER WHEEL POSITION	θ (°)	LOCK-IN READING
3	266.00	$67.2 \pm 0.3 \mu\text{V}$	1	266.00	$26.7 \pm 0.3 \mu\text{V}$
1	266.00	$2.41 \pm 0.005 \text{ V}$	1	270.00	$26.0 \pm 0.3 \mu\text{V}$
1	266.00	$1.42 \pm 0.005 \text{ V}$	1	270.10	$26.5 \pm 0.1 \mu\text{V}$
1	267.00	$0.705 \pm 0.0005 \text{ V}$	1	270.20	$41.4 \pm 0.05 \mu\text{V}$
1	267.10	$0.370 \pm 0.0005 \text{ V}$	1	270.30	$105.4 \pm 0.05 \mu\text{V}$
1	268.00	$151 \pm 0.5 \text{ mV}$	1	270.40	$277 \pm 0.5 \mu\text{V}$
1	268.20	$60.0 \pm 0.005 \text{ mV}$	1	270.50	$0.632 \pm 0.0005 \text{ mV}$
1	268.30	$18.3 \pm 0.005 \text{ mV}$	1	270.60	$1.20 \pm 0.005 \text{ mV}$
1	268.50	$0.50 \pm 0.0005 \text{ mV}$	1	270.70	$2.20 \pm 0.005 \text{ mV}$
1	268.70	$0.20 \pm 0.0005 \text{ mV}$	1	270.80	$3.87 \pm 0.005 \text{ mV}$
1	269.00	$1.02 \pm 0.0005 \text{ mV}$	1	270.90	$6.20 \pm 0.005 \text{ mV}$
1	269.30	$0.31 \pm 0.0005 \text{ mV}$	1	271.00	$0.87 \pm 0.005 \text{ mV}$
1	269.40	$1.15 \pm 0.0005 \text{ mV}$	1	271.10	$40.9 \pm 0.05 \text{ mV}$
1	269.50	$0.655 \pm 0.0005 \text{ mV}$	1	271.20	$50.1 \pm 0.05 \text{ mV}$
1	269.60	$1.71 \pm 0.3 \mu\text{V}$	1	271.30	$153 \pm 0.3 \mu\text{V}$
1	269.70	$5.92 \pm 0.3 \mu\text{V}$	1	271.40	$5.372 \pm 0.0005 \text{ V}$
1	269.80	$46.1 \pm 0.3 \mu\text{V}$	1	271.50	$0.700 \pm 0.0005 \text{ V}$

FILTER WHEEL POSITION	θ (°)	LOCK-IN READING	FILTER WHEEL POSITION	θ (°)	LOCK-IN READING
1	273.50	1.42 ± 0.005 V	3	360.00	2.78 ± 0.005 V
1	274.00	2.41 ± 0.005 V	3	361.00	2.76 ± 0.005 V
3	274.00	66.9 ± 0.05 μ V	3	362.00	2.75 ± 0.005 V
3	275.00	163 ± 0.5 μ V	3	363.00	2.73 ± 0.005 V
3	280.00	2.48 ± 0.005 mV	3	364.00	2.71 ± 0.005 V
3	280.00	38.0 ± 0.05 mV	3	365.00	2.70 ± 0.005 V
3	300.00	171 ± 0.5 mV	3	0.00	2.76 ± 0.005 V
3	310.00	0.408 ± 0.0005 V	3	1.00	2.76 ± 0.005 V
3	320.00	0.802 ± 0.0005 V	3	2.00	2.73 ± 0.005 V
3	330.00	1.56 ± 0.005 V	3	3.00	2.71 ± 0.005 V
3	340.00	2.20 ± 0.005 V	3	-1.00	2.77 ± 0.005 V
3	350.00	2.50 ± 0.005 V	3	-2.00	2.76 ± 0.005 V
3	360.00	2.77 ± 0.005 V	3	-3.00	2.77 ± 0.005 V
3	360.00	2.79 ± 0.005 V	3	0.00	2.66 ± 0.005 V
3	367.00	2.79 ± 0.005 V	1	66.50	0.980 ± 0.0005 mV
3	368.00	2.79 ± 0.005 V	1	66.60	2.56 ± 0.5 μ V
3	368.00	2.79 ± 0.005 V	1	66.70	100 ± 0.5 μ V

FILTER WHEEL POSITION	θ (°)	LOCK-IN READING	FILTER WHEEL POSITION	θ (°)	LOCK-IN READING
1	89.80	$42.2 \pm 0.05 \mu V$	1	90.500	$0.591 \pm 0.0005 mV$
1	89.90	$29.3 \pm 0.05 \mu V$			
1	89.95	$28.0 \pm 0.05 \mu V$			
1	89.96	$27.8 \pm 0.05 \mu V$			
1	89.97	$27.6 \pm 0.05 \mu V$			
1	89.98	$27.7 \pm 0.05 \mu V$			
1	89.99	$28.1 \pm 0.05 \mu V$			
1	90.00	$28.3 \pm 0.05 \mu V$			
1	90.01	$28.0 \pm 0.05 \mu V$			
1	90.02	$27.8 \pm 0.05 \mu V$			
1	90.03	$27.8 \pm 0.05 \mu V$			
1	90.04	$27.7 \pm 0.05 \mu V$			
1	90.05	$27.5 \pm 0.05 \mu V$			
1	90.100	$28.7 \pm 0.05 \mu V$			
1	90.200	$42.6 \pm 0.05 \mu V$			
1	90.300	$99.8 \pm 0.05 \mu V$			
1	90.400	$259 \pm 0.5 \mu V$			

APPENDIX B
REDUCED DATA

B	TRANSMITTANCE $\cos^4 B$	EXPERIMENTAL VALUE	NORMALIZED EXPERIMENTAL VALUE
-5.00	0.98	2.62E00 ± 0.01	0.95E00
-4.00	0.99	2.61E00 ± 0.01	0.94E00
-3.00	0.99	2.69E00 ± 0.09	0.97E00
-2.00	1.00	2.70E00 ± 0.09	0.97E00
-1.00	1.00	2.70E00 ± 0.08	0.97E00
0.00	1.00	2.70E00 ± 0.11	0.97E00
1.00	1.00	2.69E00 ± 0.07	0.97E00
2.00	1.00	2.68E00 ± 0.06	0.96E00
3.00	0.99	2.67E00 ± 0.04	0.96E00
4.00	0.99	2.63E00 ± 0.01	0.95E00
5.00	0.98	2.67E00 ± 0.01	0.96E00
10.00	0.94	2.51E00 ± 0.07	0.90E00
20.00	0.78	2.08E00 ± 0.02	0.75E00
30.00	0.56	1.51E00 ± 0.05	0.54E00
40.00	0.34	0.92E00 ± 0.02	0.33E00
50.00	0.17	0.47E00 ± 0.02	0.17E00
60.00	0.06	1.72E-01 ± 0.05E-01	0.62E-01
70.00	0.01	3.77E-02 ± 0.08E-02	1.36E-02
80.00	9.09E-04	2.48E-03 ± 0.07E-03	8.93E-04
85.00	5.77E-05	1.57E-04 ± 0.06E-04	5.65E-05
86.00	2.37E-05	6.44E-05 ± 0.2E-05	2.32E-05
86.50	1.39E-05	3.75E-05 ± 0.04E-05	1.35E-05
87.00	7.5E-06	2.02E-05 ± 0.03E-05	7.27E-06
87.50	3.62E-06	9.67E-06 ± 0.16E-06	3.48E-06
88.00	1.48E-06	3.93E-06 ± 0.03E-06	1.41E-06
88.25	8.70E-07	2.32E-06 ± 0.01E-06	8.35E-07
88.50	4.70E-07	1.26E-06 ± 0.01E-07	4.54E-07

B	TRANSMITTANCE COS ⁴ B	EXPERIMENTAL VALUE	NORMALIZED EXPERIMENTAL VALUE
88.75	2.26E-07	6.06E-07 ± 0.04E-07	2.18E-07
89.00	9.28E-08	2.50E-07 ± 0.03E-07	9.00E-08
89.25	2.94E-08	7.98E-08 ± 0.14E-08	2.87E-08
89.50	5.80E-09	1.65E-08 ± 0.06E-08	5.94E-09
89.60	2.38E-09	7.04E-09 ± 0.02E-09	2.53E-09
89.70	7.52E-10	2.75E-09 ± 0.02E-09	9.90E-10
89.75	3.62E-10	1.70E-09 ± 0.03E-09	6.12E-10
89.80	1.48E-10	1.16E-09 ± 0.03E-09	4.18E-10
89.90	9.28E-12	8.06E-10 ± 0.03E-10	2.90E-10
89.95	5.80E-13	7.70E-10 ± 0.01E-10	2.77E-10
89.96	2.38E-13	7.65E-10 ± 0.01E-10	2.75E-10
89.97	7.52E-14	7.59E-10 ± 0.01E-10	2.73E-10
89.98	1.48E-14	7.62E-10 ± 0.01E-10	2.74E-10
89.99	9.28E-16	7.73E-10 ± 0.01E-10	2.78E-10
90.00	0	7.60E-10 ± 0.45E-10	2.74E-10
90.01	9.28E-16	7.70E-10 ± 0.01E-10	2.77E-10
90.02	1.48E-14	7.65E-10 ± 0.01E-10	2.75E-10
90.03	7.52E-14	7.65E-10 ± 0.01E-10	2.75E-10
90.04	2.38E-13	7.62E-10 ± 0.01E-10	2.74E-10
90.05	5.80E-13	7.56E-10 ± 0.01E-10	2.72E-10
90.10	9.28E-12	7.89E-10 ± 0.01E-10	2.84E-10
90.20	1.48E-10	1.17E-09 ± 0.01E-09	4.21E-10
90.25	3.62E-10	1.72E-09 ± 0.06E-09	6.19E-10
90.30	7.52E-10	2.74E-09 ± 0.01E-09	9.86E-10
90.40	2.38E-09	7.12E-09 ± 0.01E-09	2.56E-09
90.50	5.80E-09	1.63E-08 ± 0.01E-08	5.87E-09
90.75	2.94E-08	7.84E-08 ± 0.05E-08	2.82E-08

B	TRANSMITTANCE COS ⁴ B	EXPERIMENTAL VALUE	NORMALIZED EXPERIMENTAL VALUE
91.00	9.28E-08	2.49E-07 ± 0.02E-07	8.96E-08
91.25	2.26E-07	6.08E-07 ± 0.03E-07	2.19E-07
91.50	4.70E-07	1.27E-06 ± 0.01E-07	4.57E-07
91.75	8.70E-07	2.36E-06 ± 0.04E-07	8.50E-07
92.00	1.48E-06	3.99E-06 ± 0.03E-06	1.44E-06
92.50	3.62E-06	9.75E-06 ± 0.07E-06	3.51E-06
93.00	7.50E-06	2.02E-05 ± 0.01E-05	7.27E-06
93.50	1.39E-05	3.74E-05 ± 0.03E-05	1.35E-05
94.00	2.37E-05	6.41E-05 ± 0.01E-05	2.31E-05
95.00	5.77E-05	1.56E-04 ± 0.01E-04	5.62E-05
100.00	9.09E-04	2.48E-03 ± 0.02E-03	8.93E-04
110.00	0.01	3.78E-02 ± 0.05E-02	1.36E-02
120.00	0.06	1.71E-01 ± 0.03E-01	6.16E-02
130.00	0.17	0.46E00 ± 0.07E-01	1.66E-01
140.00	0.34	0.96E00 ± 0.02E00	3.46E-01
150.00	0.56	1.53E00 ± 0.01E00	0.55E00
160.00	0.78	2.12E00 ± 0.05E00	0.76E00
170.00	0.94	2.59E00 ± 0.06E00	0.93E00
175.00	0.98	2.68E00 ± 0.06E00	0.96E00
176.00	0.99	2.71E00 ± 0.03E00	0.98E00
177.00	0.99	2.72E00 ± 0.01E00	0.98E00
178.00	1.00	2.73E00 ± 0.05E-01	0.98E00
179.00	1.00	2.75E00 ± 0.01E00	0.99E00
180.00	1.00	2.77E00 ± 0.03E00	1.00E00
181.00	1.00	2.76E00 ± 0.01E00	0.99E00
182.00	1.00	2.75E00 ± 0.01E00	0.99E00
183.00	0.99	2.74E00 ± 0.03E00	0.99E00

B	TRANSMITTANCE $\cos^4 B$	EXPERIMENTAL VALUE	$\log_{10} \text{TRANSMITTANCE}$	$\log_{10} (\cos^4 B)$
184.00	0.99	2.73600 ± 1.1E-03	-0.00000	-0.00000
185.00	0.98	2.73600 ± 1.1E-03	-0.00000	-0.00000
186.00	0.98	2.73600 ± 1.1E-03	-0.00000	-0.00000
187.00	0.97	2.74600 ± 1.1E-03	-0.00000	-0.00000
188.00	0.96	2.72600 ± 1.1E-03	-0.00000	-0.00000
189.00	0.95	2.69600 ± 1.1E-03	-0.00000	-0.00000
190.00	0.94	2.61600 ± 1.1E-03	-0.00000	-0.00000
200.00	0.78	2.14600 ± 1.1E-03	-0.00000	-0.00000
210.00	0.56	1.54600 ± 1.1E-03	-0.00000	-0.00000
220.00	0.34	0.94600 ± 1.1E-03	-0.00000	-0.00000
230.00	0.17	0.44600 ± 1.1E-03	-0.00000	-0.00000
240.00	0.06	1.74E-01 ± 1.1E-03	-0.00000	-0.00000
250.00	0.01	3.84E-02 ± 1.1E-03	-0.00000	-0.00000
260.00	9.09E-04	7.98E-03 ± 1.1E-03	-0.00000	-0.00000
265.00	5.77E-05	1.40E-03 ± 1.1E-03	-0.00000	-0.00000
266.00	2.37E-05	6.00E-04 ± 1.1E-03	-0.00000	-0.00000
266.50	1.39E-05	2.91E-04 ± 1.1E-03	-0.00000	-0.00000
267.00	7.50E-06	1.11E-04 ± 1.1E-03	-0.00000	-0.00000
267.50	3.62E-06	1.00E-04 ± 1.1E-03	-0.00000	-0.00000
268.00	1.48E-06	4.13E-05 ± 1.1E-03	-0.00000	-0.00000
268.25	8.70E-07	2.40E-05 ± 1.1E-03	-0.00000	-0.00000
268.50	4.70E-07	1.32E-05 ± 1.1E-03	-0.00000	-0.00000
268.75	2.26E-07	6.61E-06 ± 1.1E-03	-0.00000	-0.00000
269.00	9.28E-08	2.04E-06 ± 1.1E-03	-0.00000	-0.00000
269.10	6.09E-08	1.72E-06 ± 1.1E-03	-0.00000	-0.00000
269.20	3.80E-08	1.00E-06 ± 1.1E-03	-0.00000	-0.00000
269.25	2.94E-08	8.50E-07 ± 1.1E-03	-0.00000	-0.00000

B	TRANSMITTANCE COST \$	EXPERIMENTAL VALUE	NORMALIZED EXPERIMENTAL VALUE
269.30	2.22E-03	0.00E+0 ± 1.19E-03	0.00E+0
269.40	1.22E-03	0.00E+0 ± 1.19E-03	0.00E+0
269.50	5.60E-04	7.00E-04 ± 1.19E-03	0.00E+0
269.60	2.22E-03	7.00E-04 ± 1.19E-03	0.00E+0
269.70	7.50E-04	2.00E-04 ± 1.19E-03	0.00E+0
269.75	3.40E-04	-0.00E+0 ± 1.19E-03	0.00E+0
269.80	1.40E-03	7.00E-04 ± 1.19E-03	0.00E+0
269.90	0.20E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.00	0	7.00E-04 ± 1.19E-03	0.00E+0
270.10	0.20E-02	7.00E-04 ± 1.19E-03	0.00E+0
270.20	1.40E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.25	1.60E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.30	1.80E-03	6.00E-04 ± 1.19E-03	-0.00E+0
270.40	4.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.50	2.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.60	6.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.70	4.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.75	6.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.80	2.00E-03	7.00E-04 ± 1.19E-03	0.00E+0
270.90	0.20E-03	7.00E-04 ± 1.19E-03	0.00E+0
271.00	0.40E-03	6.00E-04 ± 1.19E-03	-0.00E+0
271.25	4.00E-02	0.00E+0 ± 1.19E-03	0.00E+0
271.30	0.70E-02	0.00E+0 ± 1.19E-03	0.00E+0
271.35	0.70E-02	0.00E+0 ± 1.19E-03	0.00E+0
271.40	1.00E-02	0.00E+0 ± 1.19E-03	0.00E+0
271.50	1.00E-02	0.00E+0 ± 1.19E-03	0.00E+0
271.60	1.00E-02	0.00E+0 ± 1.19E-03	0.00E+0

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